

1. Consider the following initial-boundary value problem for the heat equation:

$$\begin{cases} \frac{\partial u}{\partial t}(x, t) - \frac{\partial^2 u}{\partial x^2}(x, t) = 0 & \text{for } x \in (0, 1), t > 0, \\ u(x, 0) = \sin^2(2\pi x), \\ u(0, t) = u(1, t) = 0. \end{cases}$$

Find the solution u expressed as a trigonometric series.

2. Consider the following initial-boundary value problem for the Klein–Gordon equation:

$$\begin{cases} \frac{\partial^2 u}{\partial t^2}(x, t) - \frac{\partial^2 u}{\partial x^2}(x, t) - u(t, x) = 0 & \text{for } x \in (0, 1), t > 0, \\ u(x, 0) = x(1 - x), \quad \frac{\partial u}{\partial t}(x, 0) = 0, \\ u(0, t) = u(1, t) = 0. \end{cases}$$

Find the solution u expressed as a trigonometric series.

3. Consider the following initial-boundary value problem for the Schrödinger equation:

$$\begin{cases} i \frac{\partial u}{\partial t}(x, t) + \frac{\partial^2 u}{\partial x^2}(x, t) = 0 & \text{for } x \in (0, 1), t > 0, \\ u(x, 0) = \sin^2(\pi x), \\ u(0, t) = u(1, t) = 0. \end{cases}$$

Find the solution u expressed as a trigonometric series.

4. Let ψ be the solution to the following initial value problem for the biharmonic heat equation on the whole line:

$$\begin{cases} \frac{\partial \psi}{\partial t}(x, t) + \frac{\partial^4 \psi}{\partial x^4}(x, t) = 0 & \text{for } x \in \mathbb{R}, t > 0, \\ \psi(x, 0) = \psi_0(x). \end{cases}$$

Show that the solution is given by the formula

$$\psi(x, t) = \frac{1}{\sqrt{2\pi}} \frac{1}{t^{\frac{1}{4}}} \int_{-\infty}^{+\infty} G\left(\frac{x-y}{t^{\frac{1}{4}}}\right) \psi_0(y) dy,$$

where $G : \mathbb{R} \rightarrow \mathbb{R}$ is the function for which its Fourier transform is given by

$$\hat{G}(a) = e^{-a^4}.$$

(note that, a priori, the inverse Fourier transform of the above expression should be a function $G : \mathbb{R} \rightarrow \mathbb{C}$; show that G is indeed real valued, i.e. that $\overline{G(x)} = G(x)$ for any $x \in \mathbb{R}$.)

5. Let us consider the following *semi-infinite* initial-boundary value problem for the heat equation:

$$\begin{cases} \frac{\partial u}{\partial t}(x, t) - \frac{\partial^2 u}{\partial x^2}(x, t) = f(x) & \text{for } x \in (0, +\infty), t > 0, \\ u(x, 0) = 0, \\ u(0, t) = 0, \end{cases}$$

where

$$f(x) = xe^{-\frac{x^2}{2}}.$$

By extending $u(x, t)$ and $f(x)$ as *odd* functions of $x \in \mathbb{R}$, solve the above problem by applying the Fourier transform in the x -variable. Verify that the solution $u(x, t)$ that you get in this way is indeed odd and that $u(x, t)$ satisfies the required boundary condition at $x = 0$ (this should be automatically true for continuous odd functions).

6 (extra). Let $V : \mathbb{R} \rightarrow \mathbb{R}$ be a smooth function (which we will call the *potential*) and let $u : \mathbb{R} \times [0, +\infty) \rightarrow \mathbb{C}$ be a solution to the Schrödinger equation:

$$i \frac{\partial u}{\partial t}(x, t) + \frac{\partial^2 u}{\partial x^2}(x, t) - V(x)u(x, t) = 0.$$

We will assume that, at any time $t \geq 0$, we have that $u(x, t), \frac{\partial u}{\partial x}(x, t) \rightarrow 0$ as $x \rightarrow \pm\infty$.

(a) In the special case when $V(x) = 0$, find an expression for u if the initial data at $t = 0$ is given by

$$u(x, 0) = e^{i\lambda x - \frac{x^2}{2}}.$$

(b) In the general case (i.e. when V is not necessarily zero), show that the quantity

$$\int_{-\infty}^{+\infty} |u(x, t)|^2 dx$$

is constant in time (this motivates the interpretation of $|u(x, t)|^2$ as the probability density of the particle described by u). *Hint: Use the fact that $|u|^2 = \operatorname{Re}\{u \cdot \bar{u}\}$ and show first that $\partial_t|u|^2 = 2\operatorname{Re}\{\partial_t u \cdot \bar{u}\}$. Then, use the equation to reexpress $\partial_t u$ and integrate by parts in x if necessary. Note that $\partial_t(\operatorname{Re}(f)) = \operatorname{Re}(\partial_t f)$ and $\partial_t \bar{f} = \overline{\partial_t f}$.*

(c) Show that the total energy of u , defined by

$$E[u](t) = \int_{-\infty}^{+\infty} \left(\left| \frac{\partial u}{\partial x}(x, t) \right|^2 + V(x) |u(x, t)|^2 \right) dx$$

is also constant in time. *Hint: Differentiate the above expression in t like before, integrate by parts with respect to the ∂_x -derivatives and use the equation to substitute for the $\frac{\partial^2 u}{\partial x^2}$ term.*